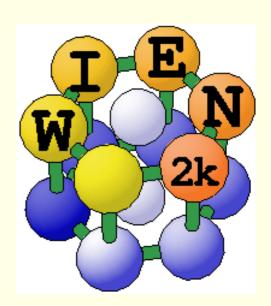




Electronic structure, atomic forces and structural relaxations by WIEN2k

Peter Blaha

Institute of Materials Chemistry
TU Vienna, Austria







- APW-based methods (history and state-of-the-art)
- WIEN2k
 - program structure + usage
 - forces, structure relaxation
- Applications
 - Phasetransitions in Aurivillius phases
 - Structure of Pyrochlore Y₂Nb₂O₇



APW based schemes



- APW (J.C.Slater 1937)
 - Non-linear eigenvalue problem
 - Computationally very demanding
- LAPW (O.K.Anderssen 1975)
 - Generalized eigenvalue problem
 - Full-potential total-energies (A.Freeman etal.)
- Local orbitals (D.J.Singh 1991)
 - treatment of semi-core states (avoids ghostbands)
- APW+lo (E.Sjöstedt, L.Nordstörm, D.J.Singh 2000)
 - Efficience of APW + convenience of LAPW
 - Basis for

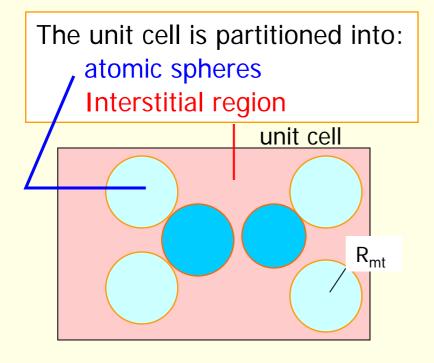


K.Schwarz, P.Blaha, G.K.H.Madsen, Comp.Phys.Commun.**147**, 71-76 (2002)

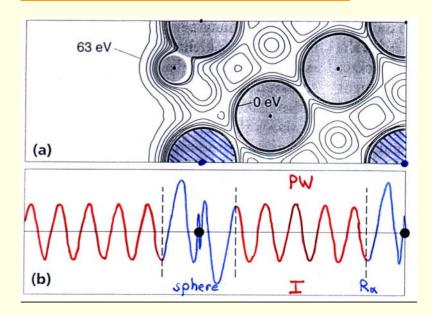


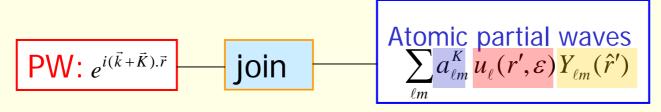
APW Augmented Plane Wave method





Bloch wave function: atomic partial waves Plane Waves (PWs)



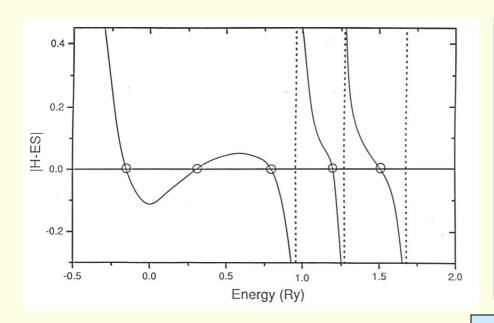


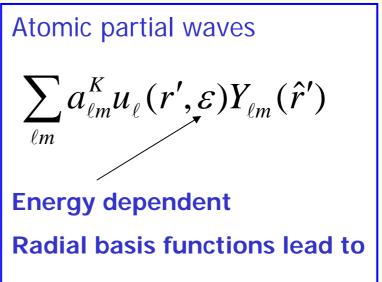
matching coefficient, radial function, spherical harmonics



Slater's APW (1937)







H Hamiltonian S Overlap matrix

Non-linear eigenvalue problem

Computationally very demanding

One had to numerically search for the energy, for which the det|H-ES| vanishes.

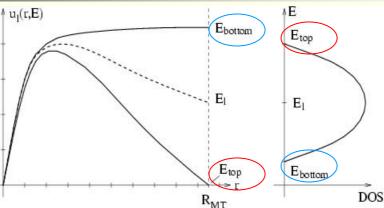


Linearization of energy dependence



LAPW suggested by

O.K.Andersen, Phys.Rev. B 12, 3060 (1975)





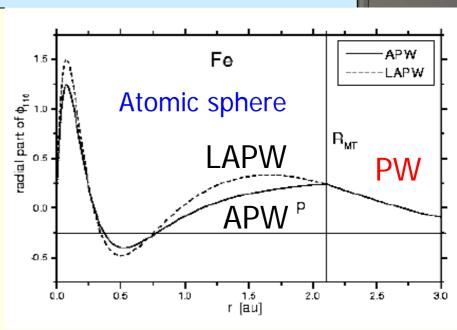
$$\Phi_{k_{n}} = \sum_{\ell m} [A_{\ell m}(k_{n})u_{\ell}(E_{\ell}, r) + B_{\ell m}(k_{n})\dot{u}_{\ell}(E_{\ell}, r)]Y_{\ell m}(\hat{r})$$

$$\dot{u}_l = \partial u_l / \partial \mathcal{E}$$
 join PWs in value and slope

 u_I expanded at E_I ("energy parameter")

→additional constraint requires more PWs than APW

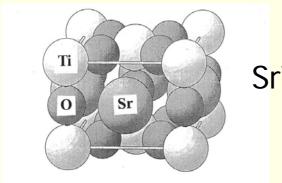
basis flexible enough for single diagonalization



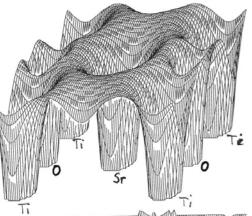


Full-potential total-energy (A.Freeman etal.)

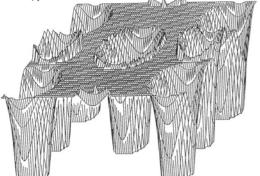




SrTiO₃



Full potential



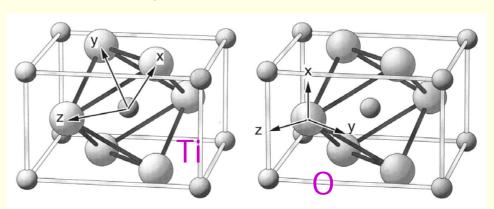
Muffin tin approximation

TiO₂ rutile

 The potential (and charge density) can be of general form (no shape approximation)

$$V(r) = \begin{cases} \sum_{LM} V_{LM}(r) Y_{LM}(\hat{r}) & r < R_{\alpha} \\ \sum_{K} V_{K} e^{i\vec{K} \cdot \vec{r}} & r \in I \end{cases}$$

 Inside each atomic sphere a local coordinate system is used (defining LM)

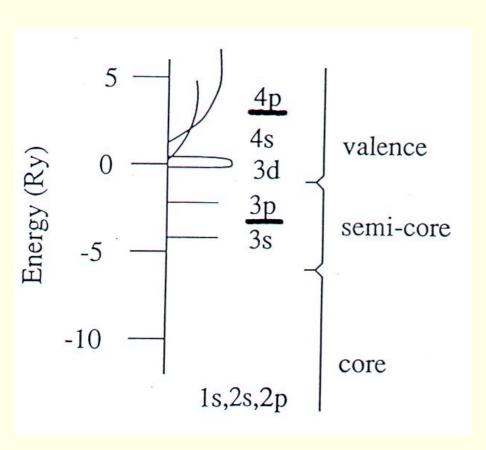




Core, semi-core and valence states



For example: Ti



Valences states

- High in energy
- Delocalized wavefunctions

Semi-core states

- Medium energy
- Principal QN one less than valence (e.g.in Ti 3p and 4p)
- not completely confined inside sphere

Core states

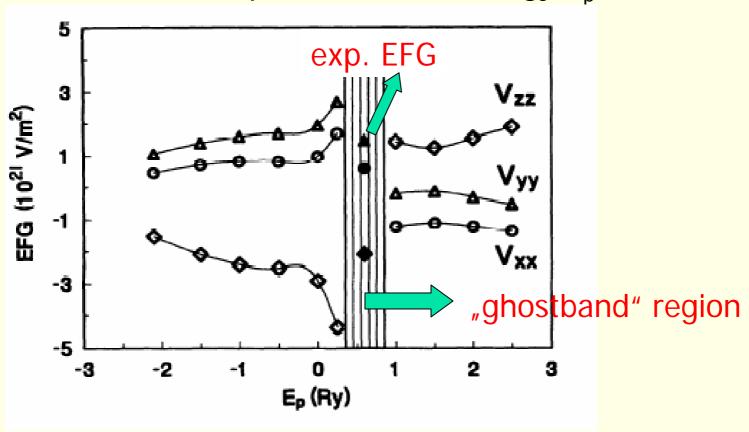
- Low in energy
- Reside inside sphere



Problems of the LAPW method:



EFG Calculation for Rutile TiO_2 as a function of the Ti-p linearization energy E_p

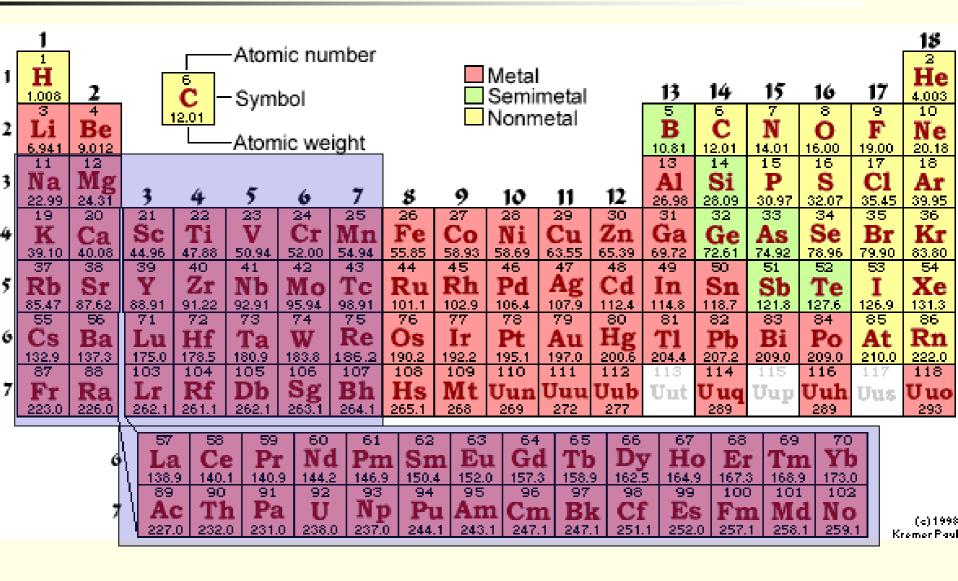


P. Blaha, D.J. Singh, P.I. Sorantin and K. Schwarz, Phys. Rev. B **46**, 1321 (1992).



Problems of the LAPW method







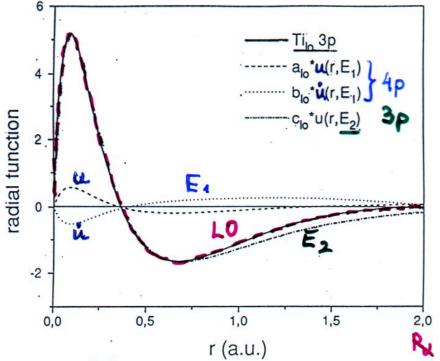
Problems with semi-core states



Extending the basis: Local orbitals (LO)







$$\Phi_{LO} = [A_{\ell m} u_{\ell}^{E_1} + B_{\ell m} \dot{u}_{\ell}^{E_1} + C_{\ell m} u_{\ell}^{E_2}] Y_{\ell m}(\hat{r})$$

- is confined to an atomic sphere
- has zero value and slope at R
- can treat two principal QN n for each azimuthal QN ℓ (3p and 4p)
- corresponding states are strictly orthogonal (no "ghostbands")
- tail of semi-core states can be represented by plane waves
- only slight increase of basis set (matrix size)

D.J.Singh, Phys.Rev. B 43 6388 (1991)



New ideas from Uppsala and Washington



E.Sjöstedt, L.Nordström, D.J.Singh, An alternative way of linearizing the augmented plane wave method, Solid State Commun. 114, 15 (2000)

- •Use APW, but at fixed E_{I} (superior PW convergence)
- Linearize with additional lo (add a few basis functions)

$$\Phi_{k_n} = \sum_{\ell m} A_{\ell m}(k_n) u_{\ell}(E_{\ell}, r) Y_{\ell m}(\hat{r})$$

$$\Phi_{lo} = [A_{\ell m} u_{\ell}^{E_1} + B_{\ell m} \dot{u}_{\ell}^{E_1}] Y_{\ell m}(\hat{r})$$

optimal solution: mixed basis (K.Schwarz, P.Blaha, G.K.H.Madsen, Comp.Phys.Commun.**147**, 71-76 (2002))

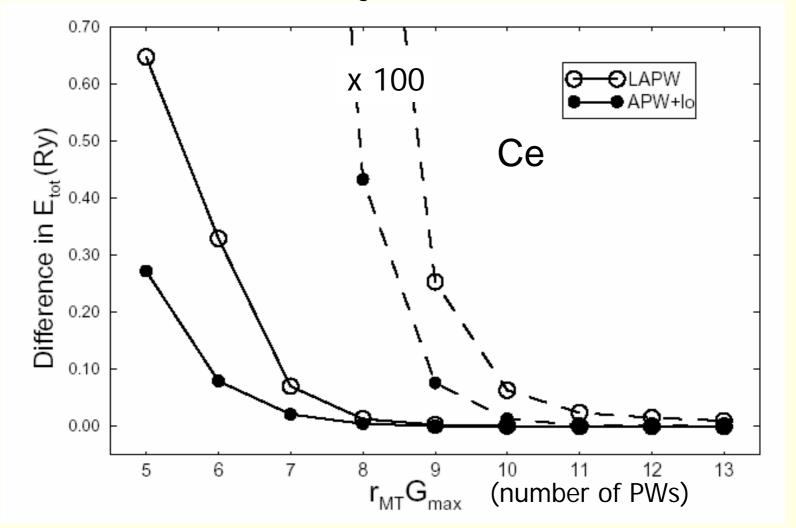
- use APW+lo for states which are difficult to converge: (f or d- states, atoms with small spheres)
- use LAPW+LO for all other atoms and angular momenta



Convergence of the APW+LO Method



E. Sjostedt, L. Nordstrom and D.J. Singh, Solid State Commun. 114, 15 (2000).





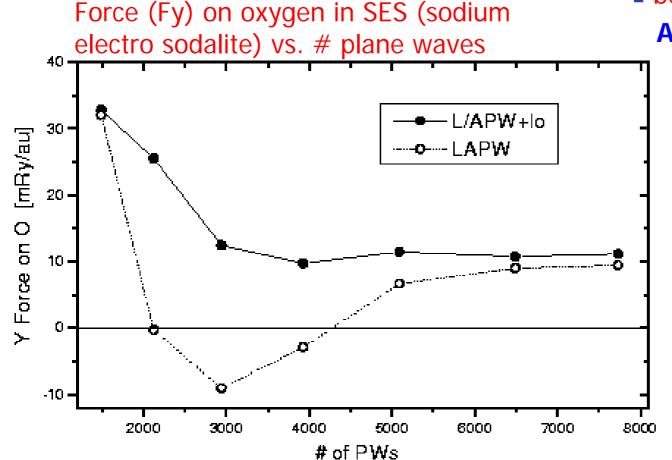
Improved convergence of APW+lo

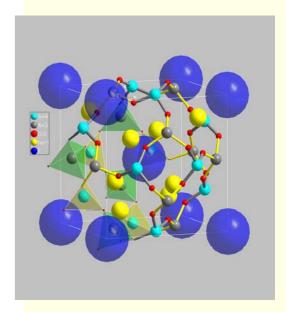


K.Schwarz, P.Blaha, G.K.H.Madsen, Comp.Phys.Commun.**147**, 71-76 (2002)

changes sign and converges slowly in LAPW

better convergence inAPW+Io

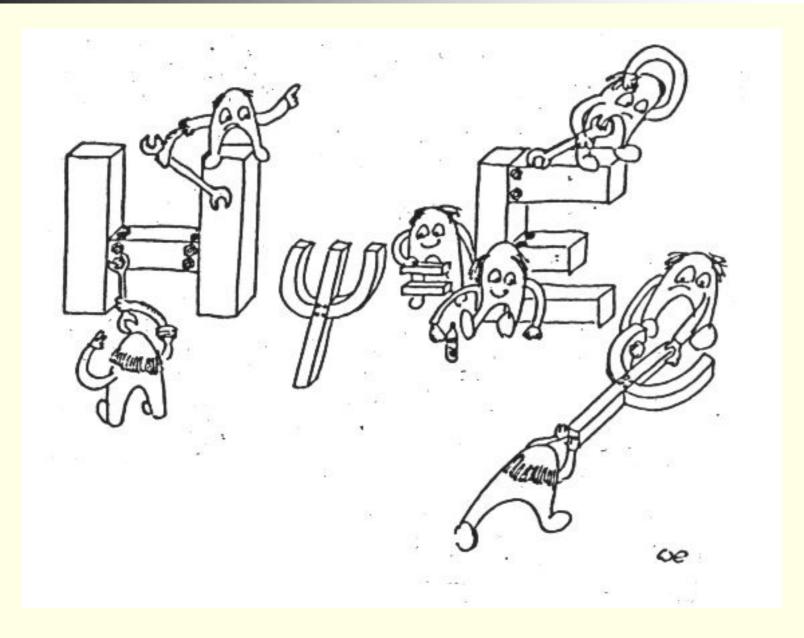






Quantum mechanics at work

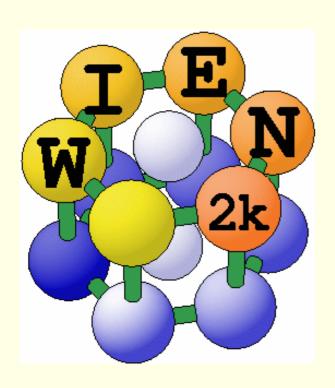






WIEN2k software package





WIEN97: ~500 users

WIEN2k: ~730 users

An Augmented Plane Wave Plus Local Orbital Program for Calculating Crystal Properties

Peter Blaha
Karlheinz Schwarz
Georg Madsen
Dieter Kvasnicka
Joachim Luitz

November 2001 Vienna, AUSTRIA Vienna University of Technology

http://www.wien2k.at



Development of WIEN2k



- Authors of WIEN2k
 - P. Blaha, K. Schwarz, D. Kvasnicka, G. Madsen and J. Luitz
- Other contributions to WIEN2k
 - C. Ambrosch-Draxl (Univ. Graz, Austria), optics
 - U. Birkenheuer (Dresden), wave function plotting
 - T. Charpin (Paris), elastic constants
 - R. Dohmen und J. Pichlmeier (RZG, Garching), parallelization
 - R. Laskowski (Vienna), non-collinear magnetism
 - P. Novák and J. Kunes (Prague), LDA+U, SO
 - B. Olejnik (Vienna), non-linear optics
 - C. Persson (Uppsala), irreducible representations
 - M. Scheffler (Fritz Haber Inst., Berlin), forces
 - D.J.Singh (NRL, Washington D.C.), local oribtals (LO), APW+lo
 - E. Sjöstedt and L Nordström (Uppsala, Sweden), APW+lo
 - J. Sofo and J. Fuhr (Barriloche), Bader analysis
 - B. Yanchitsky and A. Timoshevskii (Kiev), spacegroup
- and many others



International co-operations



More than 700 user groups worldwide

- 40 industries (Canon, Eastman, Exxon, Fuji, A.D.Little, Mitsubishi, Motorola, NEC, Norsk Hydro, Osram, Panasonic, Samsung, Sony, Sumitomo).
- Europe: (EHT Zürich, MPI Stuttgart, Dresden, FHI Berlin, DESY, TH Aachen, ESRF, Prague, Paris, Chalmers, Cambridge, Oxford)
- America: ARG, BZ, CDN, MX, USA (MIT, NIST, Berkeley, Princeton, Harvard, Argonne NL, Los Alamos Nat.Lab., Penn State, Georgia Tech, Lehigh, Chicago, SUNY, UC St.Barbara, Toronto)
- far east: AUS, China, India, JPN, Korea, Pakistan, Singapore, Taiwan (Beijing, Tokyo, Osaka, Sendai, Tsukuba, Hong Kong)

Registration at <u>www.wien2k.at</u>

- 400/4000 Euro for Universites/Industries
- code download via www (with password), updates, bug fixes, news
- usersguide, faq-page, mailing-list with help-requests



WIEN code as benchmark







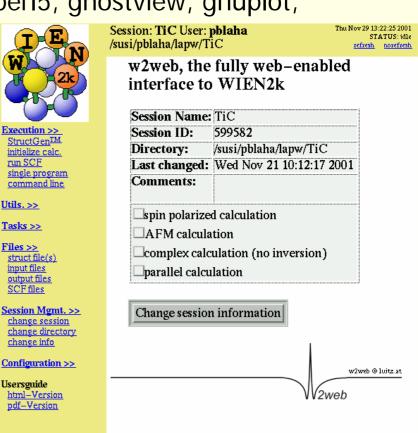
WIEN2k- hardware/software



- WIEN2k runs on any Unix/Linux platform from PCs, workstations, clusters to supercomputers
- Fortran90 (dynamical allocation)
- many individual modules, linked together with C-shell or perl-scripts
- f90 compiler, BLAS-library (mkl, Atlas), perl5, ghostview, gnuplot,

TcI/Tk (Xcrysden), www-browser

- web-based GUI w2web
- real/complex version (inversion)
- •10 atom cells on 128Mb PC
- •100 atom cells require 1-2 Gb RAM
- •k-point parallel on clusters with common NFS (slow network)
- •MPI/Scalapack parallelization for big cases (>50 atoms) and fast network
- installation support for many platforms





W2web – A web-interface

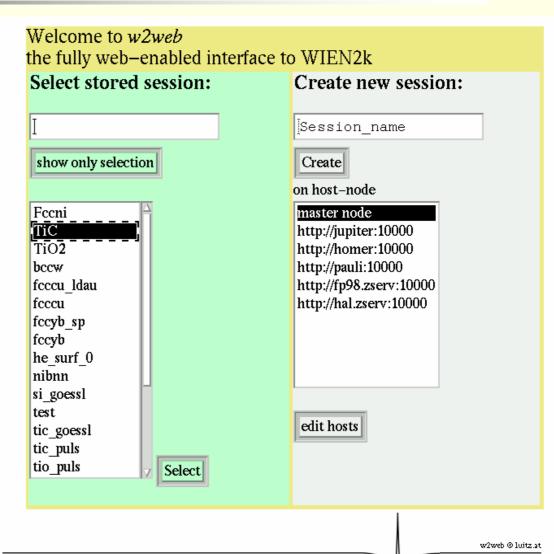


Based on www

- WIEN2k can be managed remotely via w2web
- Secure (Password)

Important steps:

- start w2web on all your hosts
- use your favorite browser and connect to the master host:port
- define a session on the desired host



V2web

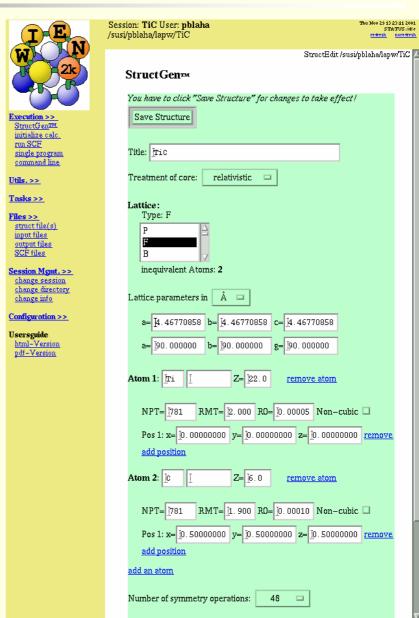


GUI (Graphical user interface)



Idea and realization by loits at @ 2001

- Structure generator
 - cif-file converter
 - spacegroup selection
- step by step initialization
 - symmetry detection
 - automatic input generation
- SCF calculations
 - Magnetism (spin-polarization)
 - Spin-orbit coupling
 - Forces (automatic geometry optimization)
- Guided Tasks
 - Energy band structure
 - DOS
 - Electron density
 - X-ray spectra
 - Optics





Program structure of WIEN2k

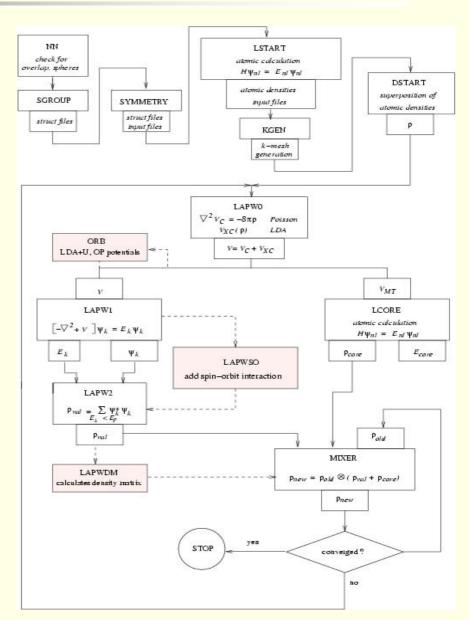


init_lapw

- initialization
- symmetry detection (F, I, Ccentering, inversion)
- input generation with recommended defaults
- quality (and computing time) depends on k-mesh and R.Kmax (determines #PW)

run_lapw

- scf-cycle
- optional with SO and/or LDA+U
- different convergence criteria (energy, charge, forces)





Task for electron density plot



STATUS:idle

refresh norefresh

- A task consists of
 - a series of steps
 - that must be executed
 - to generate a plot
- For electron density plot
 - select states (e.g. valence e-)
 - select plane for plot
 - generate 3D or contour plot with gnuplot or Xcrysden



Execution >>

StructGenIM initialize calc. run SCF single program command line

Utils.>>

<< Tasks

El. Dens.

DOS

XSPEC

ELNES

OPTIC

Bandstructure

Files >>

struct file(s) input files output files SCF files

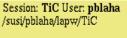
Session Mgmt. >>

change session change directory change info

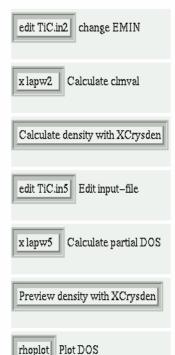
Configuration >>

Usersguide

html-Version pdf-Version

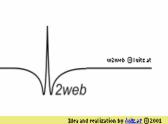


Electron density plots



rnopiot Piot Di

edit TiC.in2 reset EMIN





Properties with WIEN2k - I



Energy bands

- classification of irreducible representations
- 'character-plot' (emphasize a certain band-character)

Density of states

- including partial DOS with I and m- character (eg. p_x , p_y , p_z)
- Electron density, potential
 - total-, valence-, difference-, spin-densities, ρ of selected states
 - 1-D, 2D- and 3D-plots (Xcrysden)
 - X-ray structure factors
 - Bader ´s atom-in-molecule analysis, critical-points, atomic basins and charges $(\nabla \rho.\vec{n}=0)$
 - spin+orbital magnetic moments (spin-orbit / LDA+U)
- Hyperfine parameters
 - hyperfine fields (contact + dipolar + orbital contribution)
 - Isomer shift
 - Electric field gradients



Properties with WIEN2k - II



Spectroscopy

- core levels (with core holes)
- X-ray emission, absorption, electron-energy-loss
 - (core valence/conduction-band transitions including matrix elements and angular dep.)
 - EELS inclusion of possible non-dipol transititons (momentum transfer)
- optical properties (dielectric function, JDOS including momentum matrix elements and Kramers-Kronig)
- fermi surface (2D, 3D)

Total energy and forces

- optimization of internal coordinates, (MD, BROYDEN)
- cell parameter only via E_{tot} (no stress tensor)
- elastic constants for cubic cells
- Phonons via supercells
 - interface to PHONON (K.Parlinski) bands, DOS, thermodynamics, neutrons



Total energies and atomic forces



(Yu et al.; Kohler et al.)

Total Energy:

- Electrostatic energy
- Kinetic energy
- XC-energy

Force on atom α:

- Pulay corrections
 - Core
 - Valence
- occupied states (done only in last scf-iteration)

$$U[\rho] = \frac{1}{2} \int d^{3}\vec{r} \ \rho(\vec{r}) V_{es}(\vec{r}) + \frac{1}{2} \sum_{\alpha} Z_{\alpha} V_{es}^{\alpha}(\vec{r})$$

$$T[\rho] = \sum_{i} n_{i} \varepsilon_{i} - \int d^{3} \vec{r} \ \rho(\vec{r}) V_{eff}(\vec{r})$$

$$E_{xc}[\rho] = \int d^3 \vec{r} \ \rho(\vec{r}) \varepsilon_{xc}(\vec{r})$$

$$\vec{F}^{\alpha} = \frac{-dE_{tot}}{d\vec{R}_{\alpha}} = F_{HF}^{\alpha} + F_{core}^{\alpha} + F_{val}^{\alpha}$$

• Hellmann-Feynman-force
$$F_{HF}^{\alpha} = Z_{\alpha} \sum_{m=-1}^{1} \lim_{r_{\alpha} \to 0} \frac{V_{1m}^{es}(r_{\alpha})}{r_{\alpha}} \nabla_{\alpha} [r_{\alpha} Y_{1m}(\hat{r})]$$

$$F_{core}^{\alpha} = -\int \rho_{core}(r) \nabla_{\alpha} V_{eff}(r) \, d\vec{r}$$

$$\left[(K^2 - \varepsilon_i) \oint \phi_{K'}^*(r) \phi_K(r) \, dS_\alpha - i(K - K') \left\langle \phi_{K'} \middle| H - \varepsilon_i \middle| \phi_K \right\rangle_\alpha \right]$$

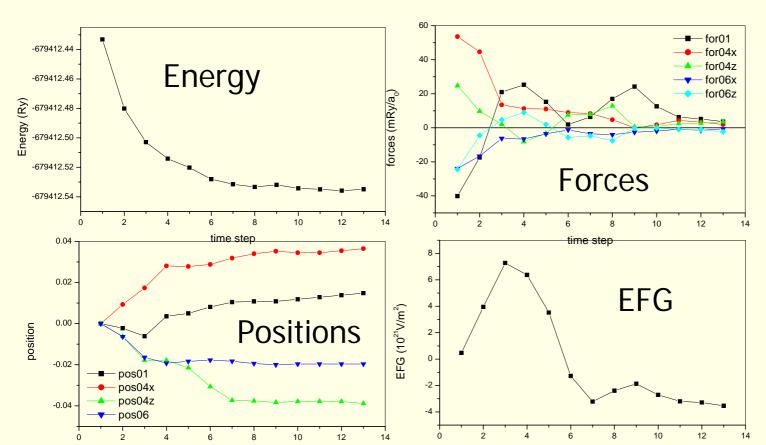


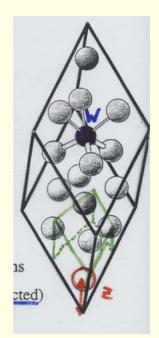
Optimization of atomic posistions (E-minimization via forces)



- damped Newton scheme
 (not user friendly, input adjustment by hand)
- new quite efficient Broyden scheme

W impurity in Bi (2x2x2 cell: Bi₁₅W)







PHONON-I



PHONON

- by K.Parlinski (Crakow)
- runs under MS-windows
- uses a "direct" method to calculate Forceconstants with the help of an ab initio program
- with these Forceconstants phonons at arbitrary k-points can be obtained
- Define your spacegroup
- Define all atoms at previously optimized positions





PHONON-II



- Define an interaction range (supercell)
 - create displacement file
 - transfer case.d45 to Unix
- Calculate forces for all required displacements
 - init_phonon_lapw
 - for each displacement a case_XX.struct file is generated in an extra directory
 - runs nn and lets you define RMT values like:
 - **1.85 1-16**



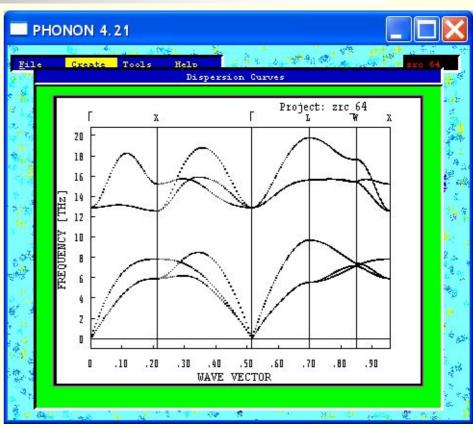
- init_lapw: either without symmetry (and then copies this setup to all case_XX)
 or with symmetry (must run init_lapw for all case_XX) (Do NOT use SGROUP)
- run_phonon: run_lapw –fc 0.1 –i 40 for each case_XX



PHONON-III



- analyze_phonon_lapw
 - reads the forces of the scf runs
 - generates "Hellman-Feynman" file case.dat and a "symmetrized HFfile case.dsy (when you have displacements in both directions)
 - check quality of forces:
 - sum F_x should be small (0)
 - abs(F_x) should be similar for +/displacements
- transfer case.dat (dsy) to Windows
- Import HF files to PHONON
- Calculate phonons





Ferroelectricity in Aurivillius phases

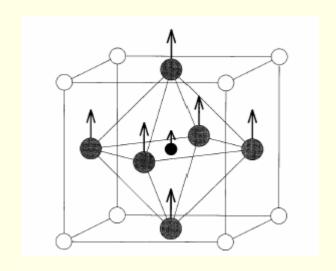


- *Bilbao: J. Manuel Perez-Mato, M. Aroyo
- Universidad del País Vasco
- *Vienna: P. Blaha, K. Schwarz,
- J. Schweifer
- *Cracow: K. Parlinski



Perovskites: PbTiO₃





Classical Example of a Ferroelectric

PbTiO₃

(perovskite distorted structure)

Features:

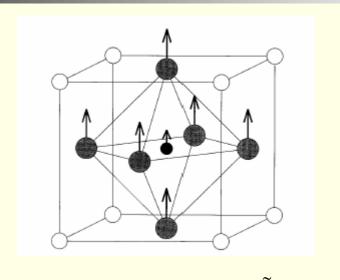
- spontaneous polarization- frozen polar mode
- phase transition:

paraelectric ----- ferroelectric



Perovskites: PbTiO3





PbTiO₃

• Symmetry break at Tc:

cubic ---- tetragonal Pm-3m ----- P4mm

- A single (degenerate) *normal* mode responsible for the transition
- Structure in Ferroelectric Phase:

high-symmetry structure + frozen polar mode

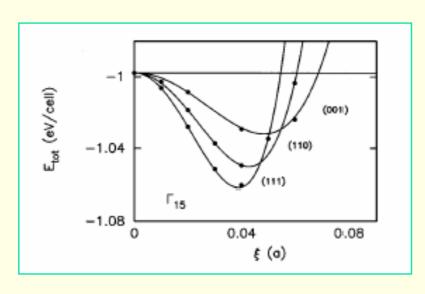
• Strain: spontaneous strain $\Delta \epsilon - 10^{-2}$



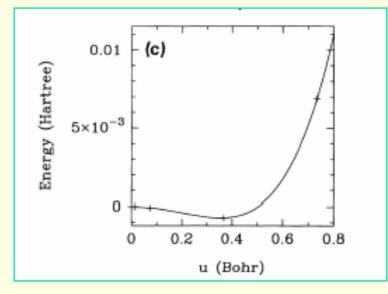
Ab-initio "prediction" of the ferroelectric inestability



PbTiO₃



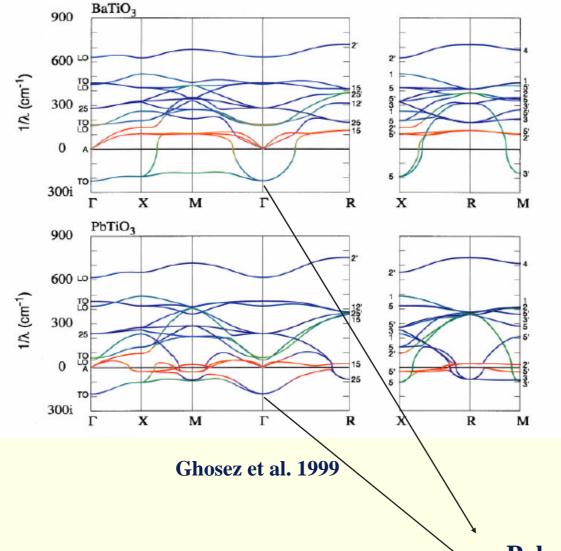
Waghmare&Rabe 1997



King-Smith&Vanderbilt 1994







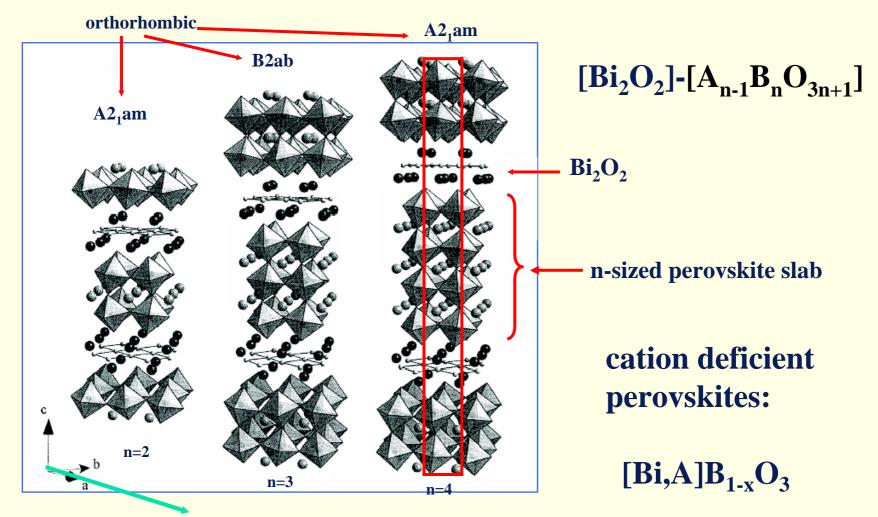
Ab-initio Phonon Branches

Polar unstable mode



The Aurivillius Compounds

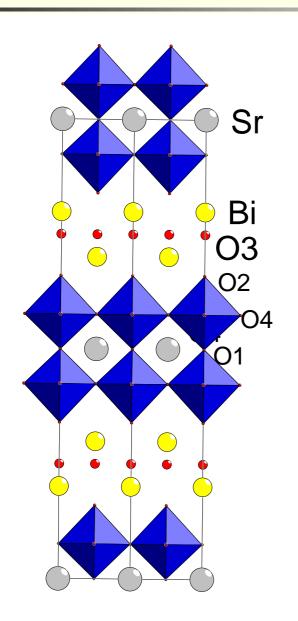




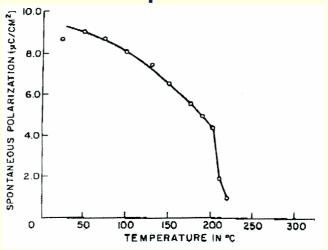


$Bi_2SrTa_2O_9$ - SBT

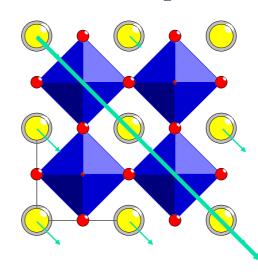




Spontaneous Polarization



polar mode Eu (deg. 2)

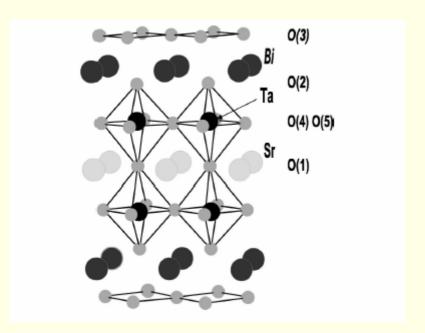


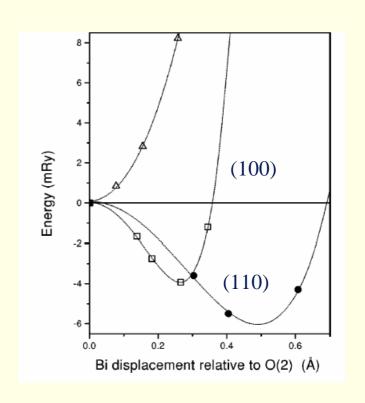


Previous ab-initio calculations in SBT



Stachiotti et al. 2000



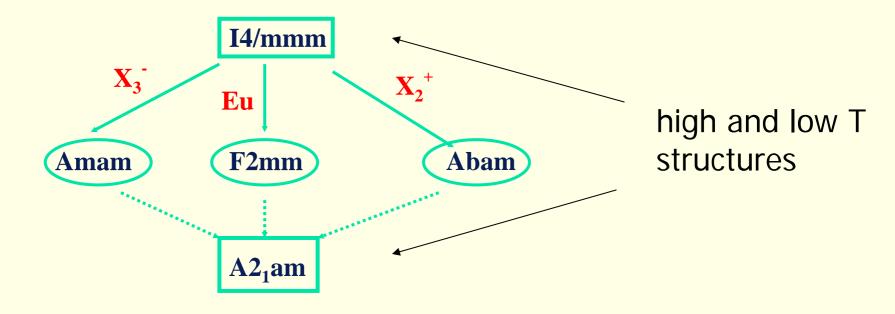


- Unstable Eu polar mode
- Strong contribution of Bi displacement : Bi-O(2) hybridization



SBT – An Usual Ferroelectric? NO!





From symmetry analysis (group/ subgroup relations) the Eu mode alone cannot explain the complete phase transition

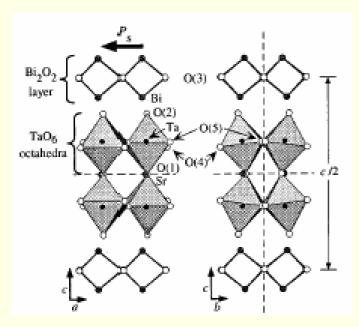
Ferroelectric Phase = mode Eu + mode X3- + mode X2+



Frozen distortions in SBT

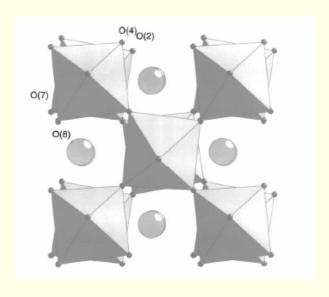






Shimakawa et al. 1999

mode X₂⁺



Hervoches et al. 2002

Ferroelectric Phase = mode Eu + mode X3 - + mode X2 +

Eigenvectors:

e_{Eu}

 e_{X3}

 $\mathbf{e}_{\mathbf{X2}}$

Mode Amplitudes (bohrs) : QEu = 0.82

QX3 = -1.20

QX2 = -0.35



All Aurivillius: "Atypical" Ferroelectrics:



Prototype tetragonal paraelectric phase

Not possible through the freezing of a polar mode!

Ferroelectric orthorhombic phase

•A single order parameter can not produce this transition.

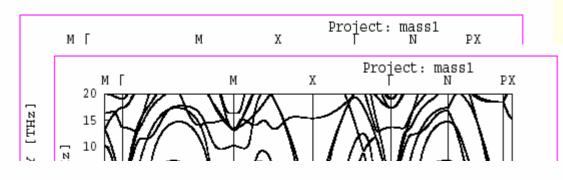
According to Landau theory:

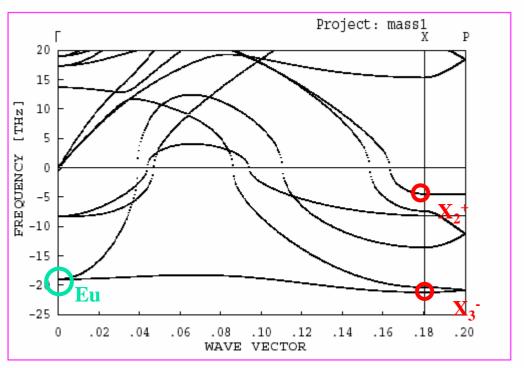
•Either discontinuous transition or intermediate phases



Phonons in SBT

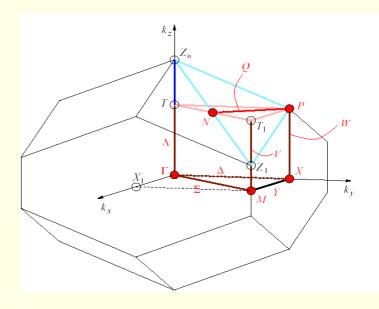






SBT

CALCULATED "PHONON" BRANCHES



Comparing calculated and exp. soft-modes Eigenvectors:



Ferroelectric Phase = mode Eu + mode X3- + mode X2+

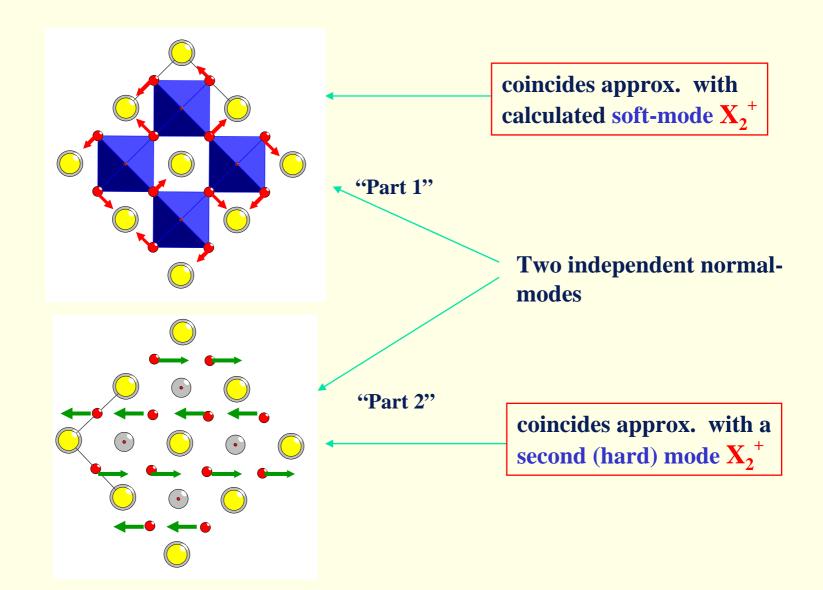
In general, good agreement with experimental frozen modes

Except !.... The experimental frozen mode X_2^+



Comparing exp. and calculated X_2^+ Modes

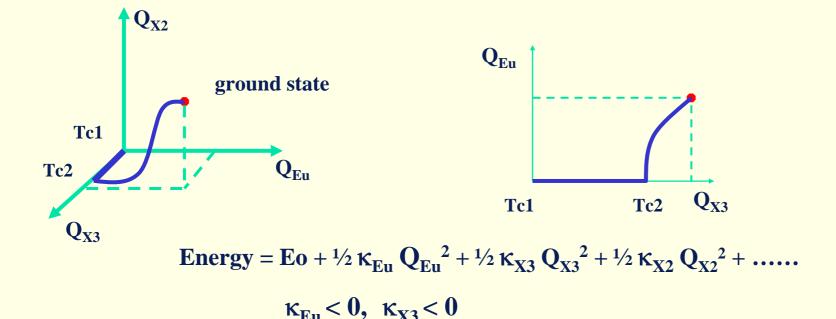






Energy Landscape around the tetragonal configuration





Transition Sequence:

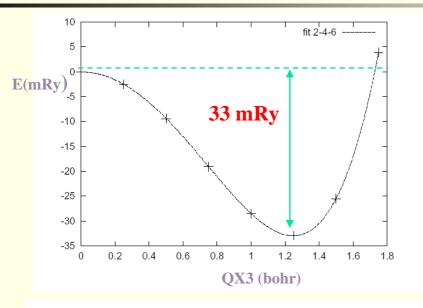
Thermal Free Energy = Fo + $\frac{1}{2} \alpha_{Eu}(T) Q_{Eu}^2 + \frac{1}{2} \alpha_{X3}(T) Q_{X3}^2 + \frac{1}{2} \alpha_{X2} Q_{X2}^2 + \dots$



ENERGY MAP – SBT

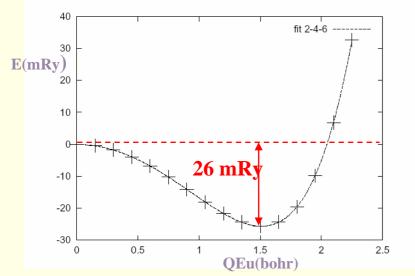
Main axes







$$E_{X3} = -40.2 Q_{X3}^2 + 10.9 Q_{X3}^4 + 0.87 Q_{X3}^6$$



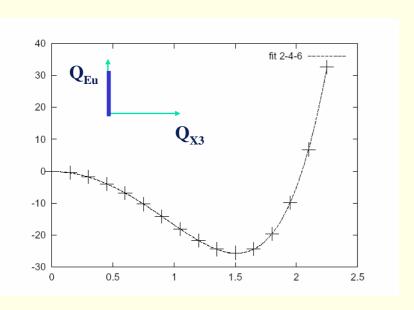


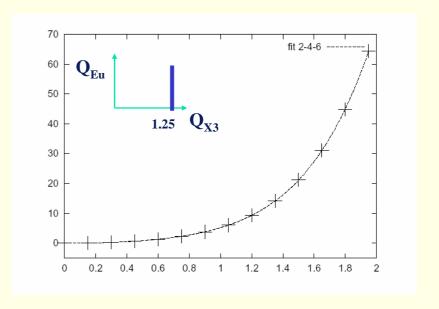
$$E_{Eu} = -19.9 Q_{Eu}^2 + 2.6 Q_{Eu}^4 + 0.52 Q_{Eu}^6$$



ENERGY MAP – SBT: Coupling Eu $-X_3^-$







$$E_{Eu} = -19.9 Q_{Eu}^2 + 2.6 Q_{Eu}^4 + 0.51 Q_{Eu}^6$$

$$E_{Eu} = +3.1 Q_{Eu}^2 + 1.6 Q_{Eu}^4 + 0.53 Q_{Eu}^6$$

 $+14.7 Q_{Eu}^{2} Q_{X3}^{2} - 0.63 Q_{Eu}^{4} Q_{X3}^{2}$

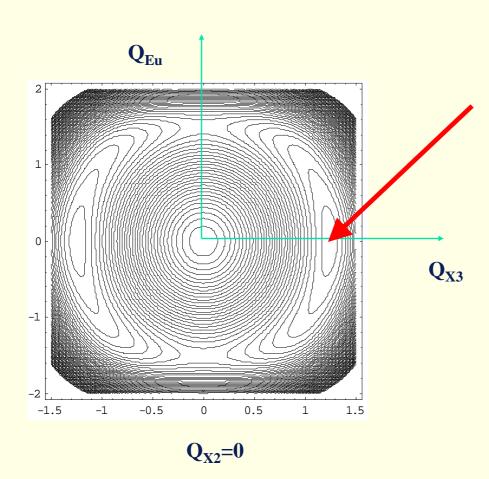
Strong biquadratic coupling Eu –X₃



SBT – Energy Map:

Coupling Eu –X₃-





Predicted Ground State:

Only mode X₃⁻ frozen!

Space group Amam

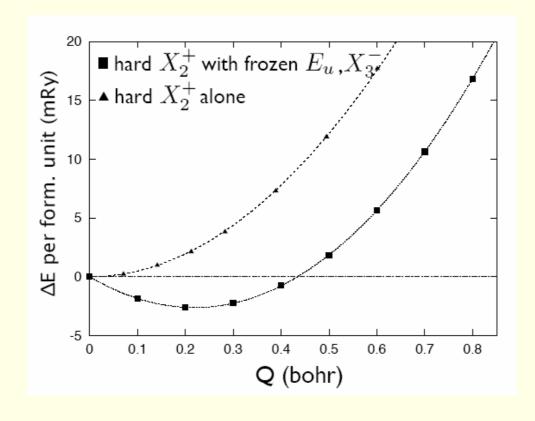
Essential reason: Strong biquadratic coupling penalizes "mixed" states

Something missing? •coupling with X₂⁺

Coupling with X₂⁺



- Coupling with soft X₂⁺ mode relatively weak !
- Strong renormalization of hard (second) X₂⁺ mode

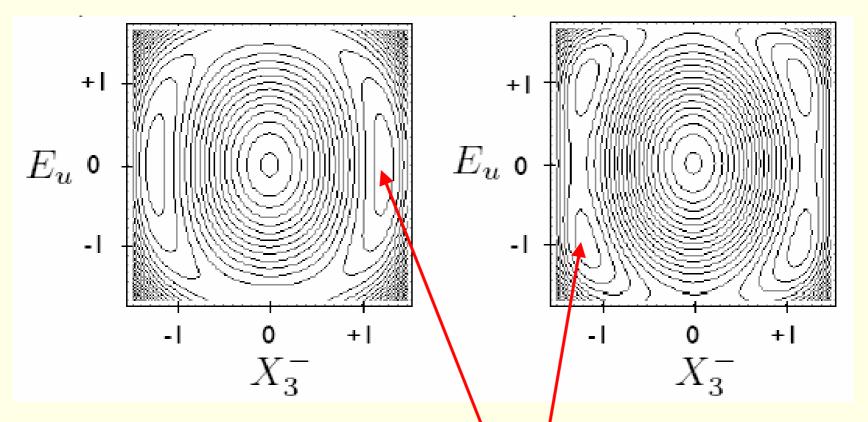




Energy landscape in SBT



including only soft X₂⁺ mode with additional hard X₂⁺ mode



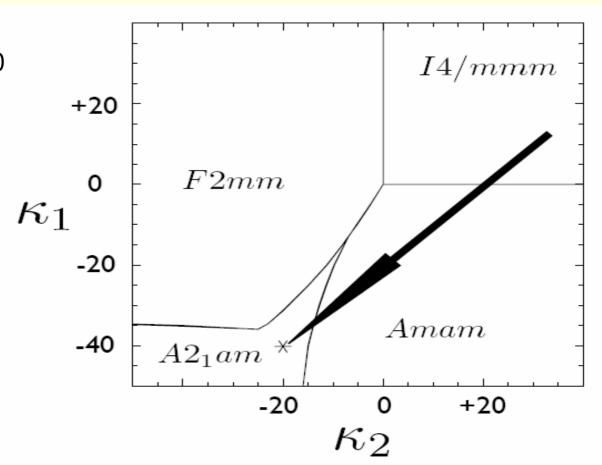
only with the additional hard X₂+\model modes can be established



Phase diagram for Bi₂SrTa₂O₉ - SBT



- Finite temperature renormalization of the T=0 energy map
- Landau theory suggests linear T-variation of the quadratic stiffness coefficients κ_{Eu} and κ_{X3}^{-}



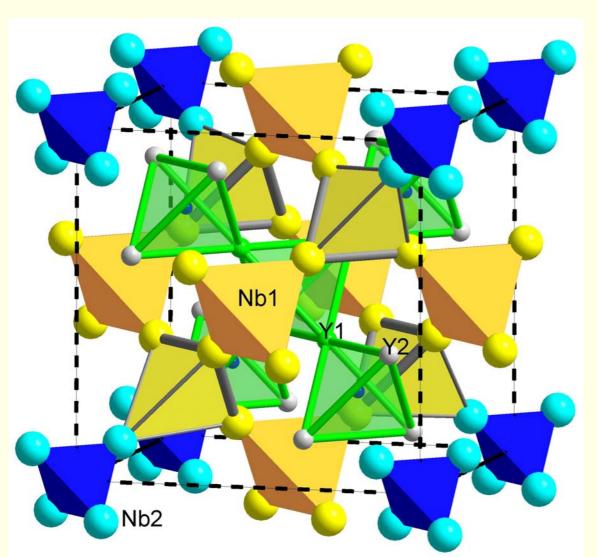
The topolgy of this phase diagram predicts **two second order** phase transitions and implies the existence of an intermediate phase of Amam symmetry. A **single first order** PT is impossible.



Pyrochlore Y₂Nb₂O₇



Insulating and non-magnetic 4d TM-oxide



P.Blaha, D.Singh, K.Schwarz, PRL in print (2004)







Metal Sublattice: Corner-shared tetrahedral network



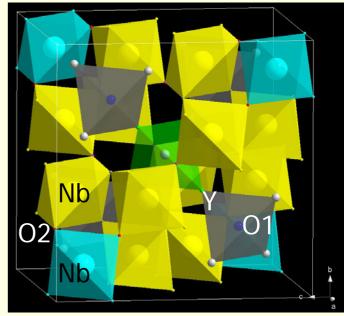
Pyrochlore structure of Y₂Nb₂O₇:

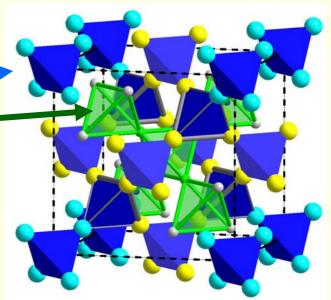


Experimental structure from powder diffraction (Fukazawa, Maeno):

Y.Maeno et al, Nature 372, 532 (1994) H.Fukazawa, Y.Maeno, Phys. Rev. B 67, 054410 (2003)

- Pyrochlore structure: Fd-3m, a_0 =10.33 Å
- 22 atoms (2 FU/cell)
- One free parameter for O2: (0.344,1/8,1/8)
- Corner shared Nb O octahedra
- Corner shared O1 Y tetrahedra
- Y surrounded by 8 oxygens (~ cube)
- Metal sublattices
 - Nb-tetrahedra (empty)
 - Y-tetrahedra (around O)
- Insulator
- Nonmagnetic





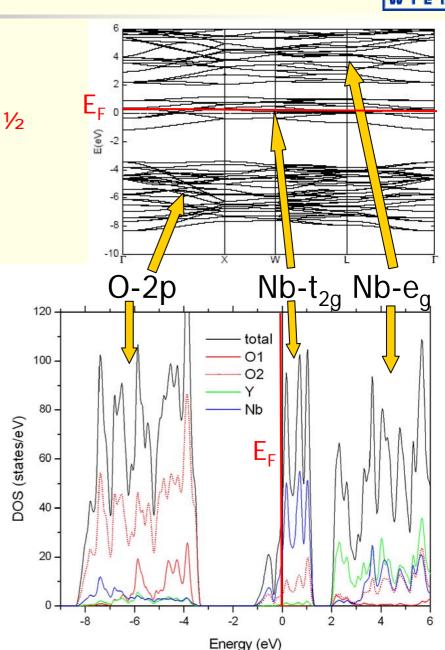


First theoretical results:



- Ionic model: Y₂³⁺Nb₂⁴⁺O₇²⁻ Nb⁴⁺: 4d¹
- → metallic or localized system with spin ½ (neither one observed in exp.)
- LDA gives nonmagnetic metallic ground-state with conventional t_{2g}-e_g splitting due to the octahedral crystal field of the oxygen atoms.

"degenerate" t_{2g} states are only partly filled.





How could one obtain a non-magnetic insulator?

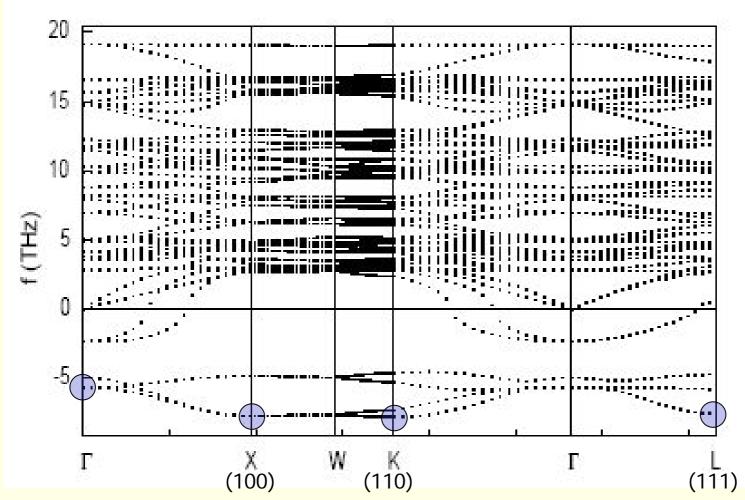


- Antiferromagnetic s=1/2 solution
 - (on geometrically frustrated lattice !?)
- Localization, strong e⁻-e⁻ correlation:
 - 4d (not 3d!) electrons,
 - thus correlation should be small (Hubbard-U ~ 2-3 eV)
 - LDA+U with U=6 eV gives insulator (FM ground state, no AFM)
 - bandwidth of t_{2q} bands: 2.5 eV (similar to U)
- structural distortion, which breaks the dominant octahedral crystal field
- Search for phonon-instabilities
 - 88 atom supercell, 46 symmetry adapted selected distortions from PHONON; resulting forces → back into PHONON



Phonon-bandstructure of Y₂Nb₂O₇





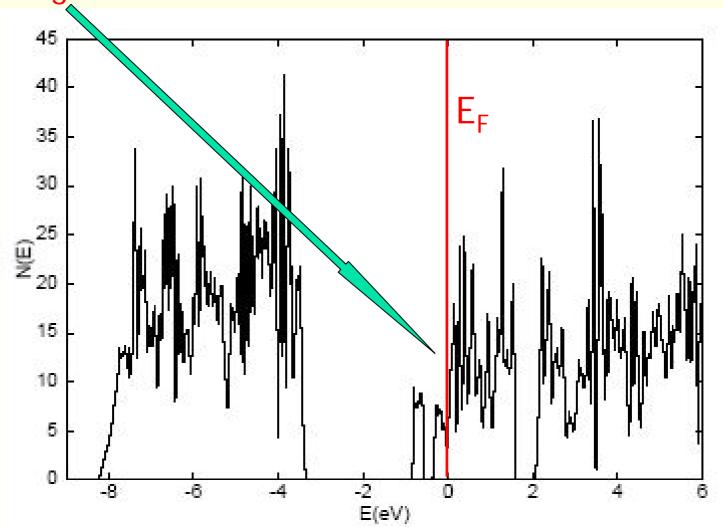
- strong Phonon-instabilities, lowest at X, K, L
- select a certain (unstable) phonon, freeze it with certain amplitude into the structure and perform full structural optimization



Γ , X and K-point phonons:



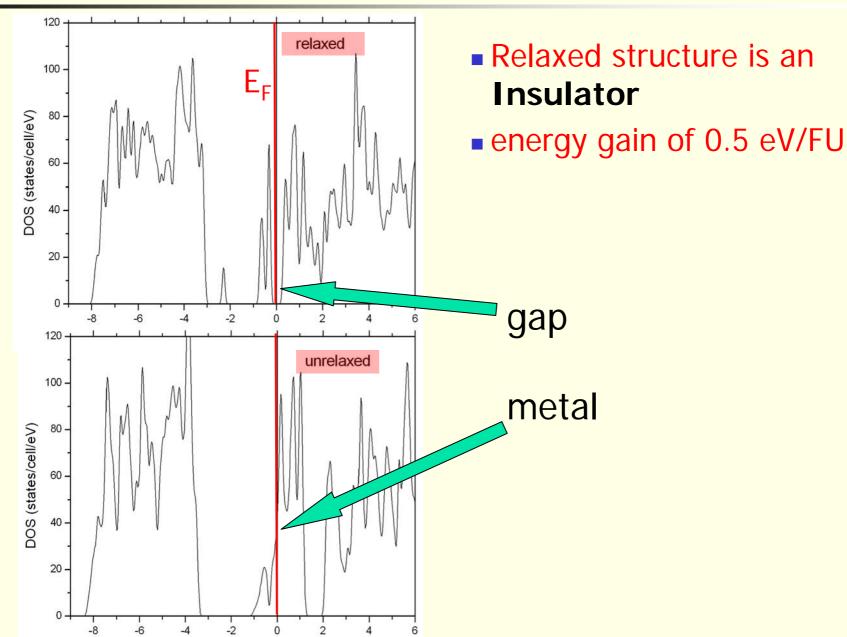
energy lower than in ideal pyrochlore structure, but still not insulating





L-point (111) phonon:







Relaxed structure:

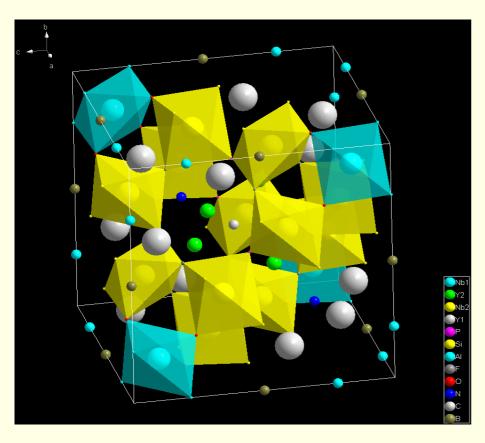


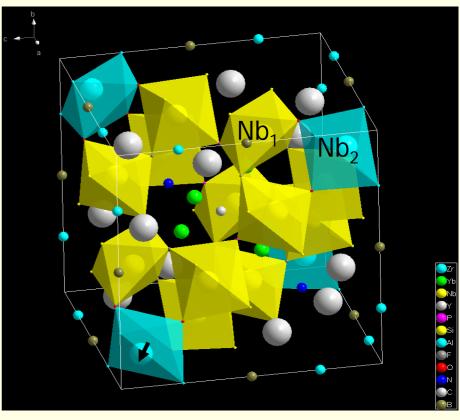
- Primitive supercell with 88 atoms
- all atoms inequivalent due to numerical optimization of the positions in P1
- Symmetrization using KPLOT (R.Hundt, J.C.Schön, A.Hannemann, M.Jansen: Determination of symmetries and idealized cell parameters for simulated structures, J.Appl.Cryst. 32, 413-416 (1999))
 - Tests possible symmetries with increasing tolerance
- Space group → P-43m, 88 atoms/cell,
 - Inequivalent atoms:
 - 2 Y
 - 2 Nb
 - **3** 01
 - **5** 02



Relaxed structure







Experimental (ideal)

relaxed:

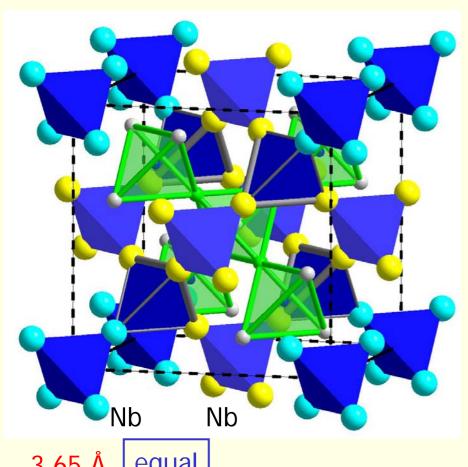
Nb₂ moves "off-center", O octahedra changes little, Nb₂-Nb₂ bonds?



Main change in structural relaxation

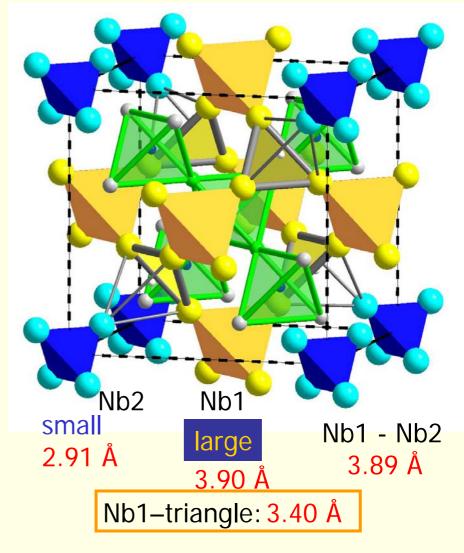


Original pyrochlore



3.65 Å equal

Relaxed structure

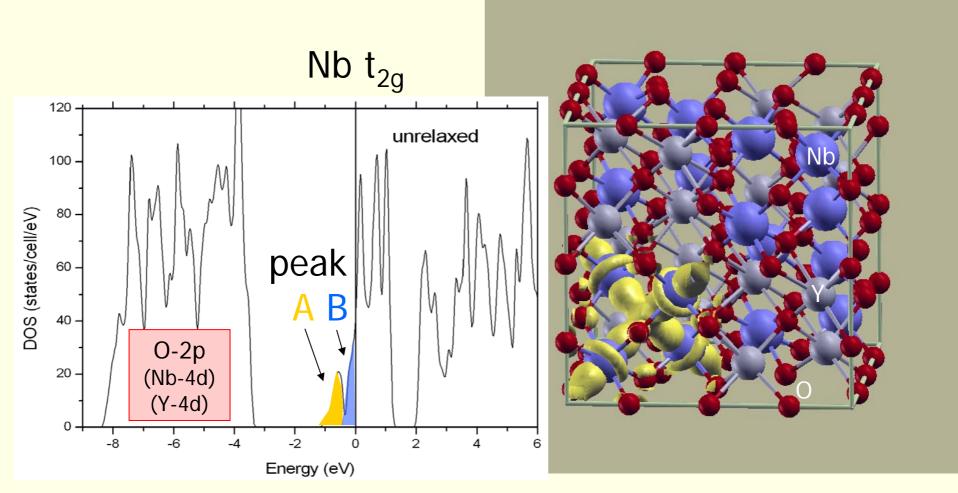




Chemical bonding: original structure



peak A: Nb-Nb bonds within the Nb tetrahedra by d-z² orbitals

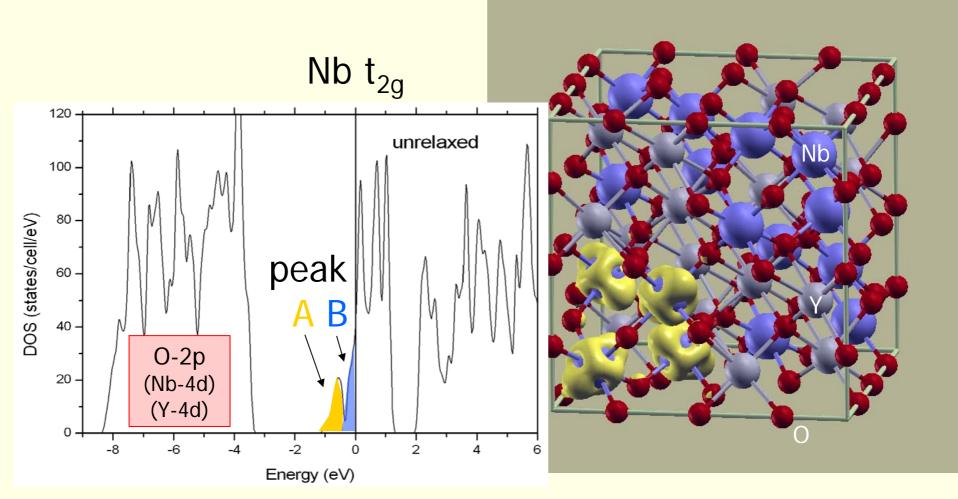




Chemical bonding: original structure



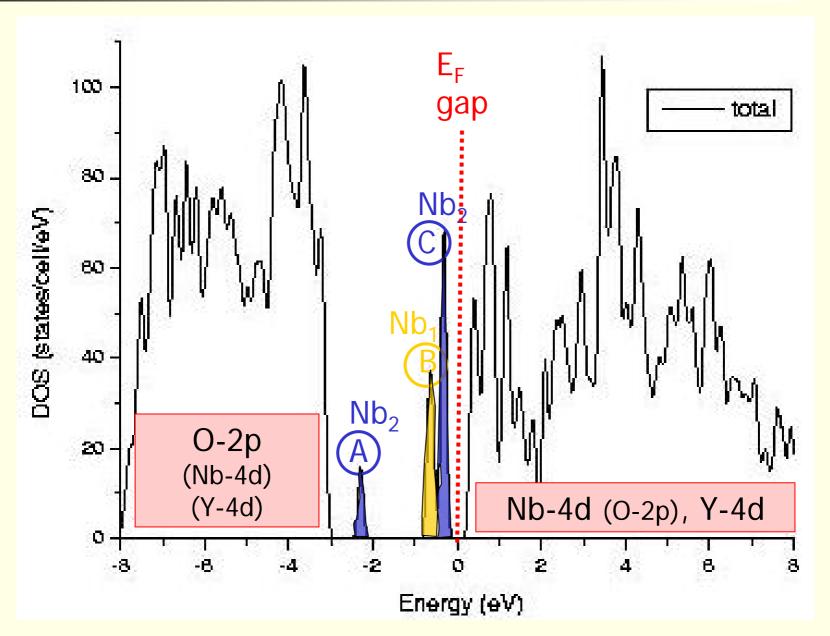
peak B: mixture of all 3 t_{2g} orbitals





DOS of relaxed structure

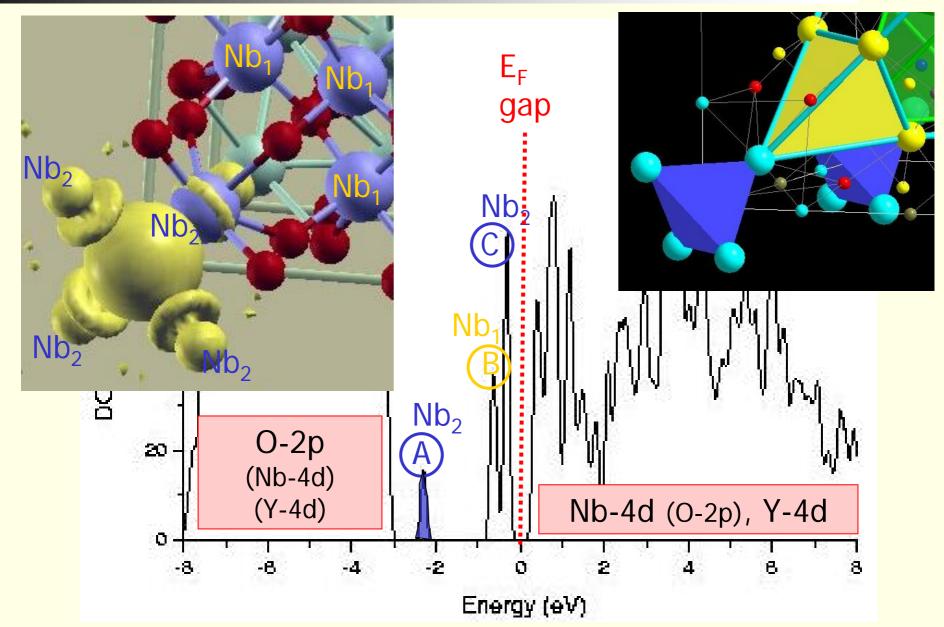






Peak A (Nb₂) 4-center bond (d-z²)

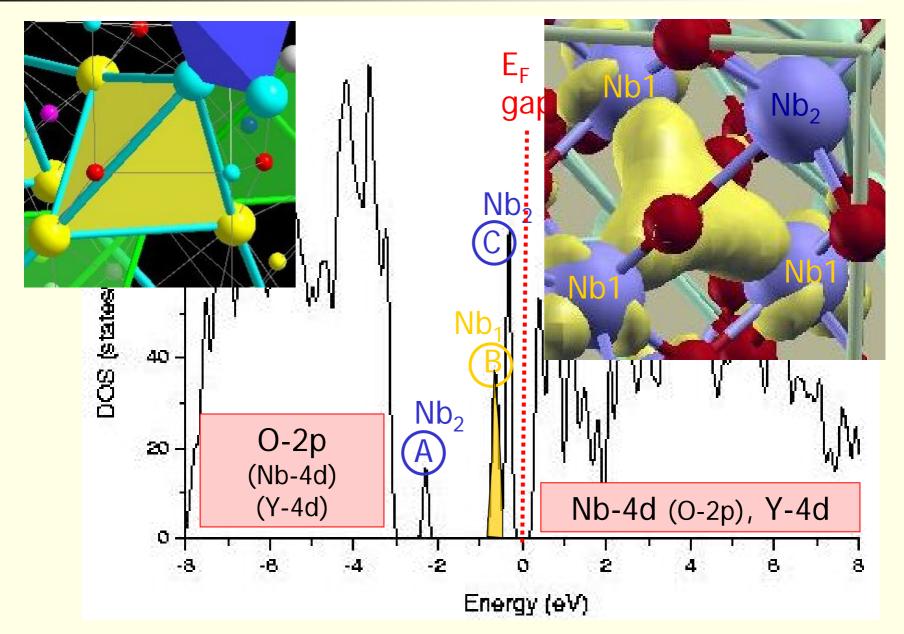






Peak B (Nb₁) 3-center bond (new !)

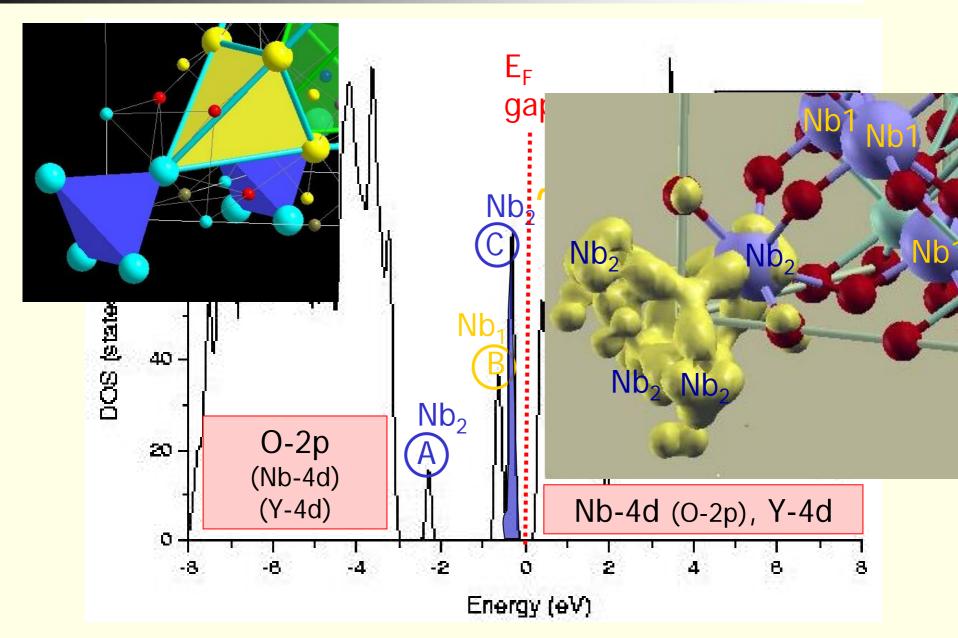






Peak C (Nb₂) 3-center bonds an all 4 faces







Results for Y₂Nb₂O₇



- Orbital ordering of Nb-4d orbitals and structural distortion!
 - Can be described with LDA !!
- M-M interaction (second NN): 2 types of bonds:
 - strong Nb-Nb 4-center bonds (only Nb2 d_z²)
 - strong Nb-Nb 3-center bonds (both, on Nb1 and Nb2)

Relaxation

- Two types of Nb
 - Nb1 (12i) highly elongated tetradedron along (111)
 - Nb2 (4e) perfect tetrahedron with short distances

Metal/insulator

Ideal structure metallic, relaxed structure insulating

P.Blaha, D.J.Singh, K.Schwarz, Charge singlets in pyrochlore Y₂Nb₂O₇ Phys.Rev.Lett. (in print) (2004)



Advantage/disadvantage of WIEN2k



- + robust all-electron full-potential method
- + unbiased basisset, one convergence parameter (LDA-limit)
- + all elements of periodic table (equal expensive), metals
- + LDA, GGA, meta-GGA, LDA+U, spin-orbit
- + many properties
- + w2web (for novice users)
- ? speed (+ memory requirements)
 - + very efficient basis for large spheres (2 bohr) (Fe: 12Ry, O: 9Ry)
 - less efficient for small spheres (1 bohr) (O: 25 Ry)
 - large cells (n³, iterative diagonalization not perfect)
- no stress tensor
- no linear response





Thank you for your attention!